kT factorisation and quarkonium production in the LHC era



Stéphane Delorme (on behalf of the 12-147 collaboration)

COPIN-IN2P3 Workshop 2023



Warsaw University of Technology



THE HENRYK NIEWODNICZAŃSKI INSTITUTE OF NUCLEAR PHYSICS POLISH ACADEMY OF SCIENCES



Laboratoire de Physique des 2 Infinis

Scientists in the collaboration

- IN2P3 scientists:
 - J.P. Lansberg (IJCLAB)
 - S. Wallon (IJCLAB)
 - M. Nefedov (IJCLAB, post-doc)
 - C. Flett (IJCLAB, post-doc)
 - S. Nabeebaccus (IJCLAB, post-doc)
 - Y. Yedelkina (IJCLAB, Ph. D student)
 - E. Li (IJCLAB, Ph. D student)
 - J. Bor (IJCLAB, Ph. D student)

- COPIN scientists:
 - A. Kusina (IFJ PAN)
 - D. Kikoła (WUT)
 - J. Wagner (NCBJ)
 - L. Szymanowski (NCBJ)
 - A. van Hameren (IFJ PAN)
 - S. Delorme (IFJPAN, post-doc)
 - A. Colpani Serri (WUT, Ph. D student)
 - A. Safronov (WUT, Ph. D student)
 - L. Manna (WUT, Ph. D student)

Exchanges in 2023

- 0 days approved for IN2P3 scientists in COPIN
- 34 days approved for COPIN scientists in France
 - A. Colpani Serri (7 days)
 - S. Delorme (7 days)
 - K. Kutak (4 days)
 - A. van Hameren (4 days)
 - T. Altinoluk (4 days)
 - G. Beuf (4 days)
 - L. Manna (4 days)
- A. Colpani Serri stayed for 3 weeks to attend doctoral school lectures in IJCLab
- L. Manna stayed for 3 weeks in September in IJCLab
- J. Wagner stayed for 2 weeks in Ecole Polytechnique
- L. Szymanowski stayed for several weeks in France

Workshops/Meetings

- Quarkonia As Tools 2023, 4–14 Jan 2023, Aussois, France
- ▶ QCD Evolution Workshop 2023, Orsay, France, 22–26 May 2023
- ▶ NLOAccess/Precisonium annual meeting, 20 24 Nov 2023, CERN
- ▶ Quarkonia As Tools 2024, 7–13 Jan 2024, Aussois, France
- Workshop on overlap between QCD resummations 14 Jan 17 Jan 2024, Aussois, France
- Usually from 5 to 15 participants from COPIN/IN2P3 (e.g. for QaT2024: S. Delorme, L. Manna, A. Colpani Serri, A. Safronov, A. Matyja on the polish side)

Ph.D students

- A. Colpani Serri: Works in WUT under the supervision of D. Kikoła (+ partial co-supervision of J.P. Lansberg) on the extension of Madgraph to TMD factorisation and to quarkonium production. Also works with C. Flett and J. Bor.
- A. Safronov: Works in WUT under the supervision of D. Kikoła (+ partial co-supervision of J.P. Lansberg) on the implementation of the computation of cross sections in proton-nucleus collisions in Madgraph at NLO.
- L. Manna: Works in WUT under the supervision of D. Kikoła (+ partial co-supervision of J.P. Lansberg) on the implementartion of the computation of cross sections in photon-hadron and lepton-hadron collisions in Madgraph at NLO.
- Strongly consolidate the contribution of WUT in NLOAccess (EU-funded Virtual Access for automated pQCD computations).

Ph.D students

- Y. Yedelkina: Works in IJCLAB on the quarkonium production with QCD corrections. Works with M. Nefedov
- E. Li: Works in IJCLAB under the supervision of S. Wallon and L. Szymanowski on probing gluon saturation in semi-hard γ(*) + p/A processes. Defended her Ph. D on 19/10/2023.
- ▶ J. Bor: Connected to COPIN-IN2P3, works on TMD factorization with A. Colpani Serri
- ► K. Lynch: Also connected to the collaboration, works on inclusive photoproduction at the LHC in UPC (related to the work of A. Safronov, L. Manna and A. Colpani Serri)

Problem: Perturbative instability of quarkonium total cross sections **Inclusive** η_c -hadroproduction [Mangano et.al., '97,

..., Lansberg, Ozcelik, '20]

At sufficiently high energy, the $\mu_{F,R}$ -variation band of NLO 'computation blows-up:

$$p+p \to c\bar{c} \begin{bmatrix} {}^{1}S_{0}^{[1]} \end{bmatrix} + X, \text{ LO: } g(p_{1}) + g(p_{2}) \to c\bar{c} \begin{bmatrix} {}^{1}S_{0}^{[1]} \end{bmatrix}$$
$$\sigma(\sqrt{s_{pp}}) = f_{i}(x_{1}, \mu_{F}) \otimes f_{j}(x_{2}, \mu_{F}) \otimes \hat{\sigma}(z),$$

where $z = \frac{M^2}{\hat{s}}$ with $\hat{s} = (p_1 + p_2)^2$.



$$\gamma + p \rightarrow c\bar{c} \begin{bmatrix} {}^{3}S_{1}^{[1]} \end{bmatrix} + X, \text{ LO: } \gamma(q) + g(p_{1}) \rightarrow c\bar{c} \begin{bmatrix} {}^{3}S_{1}^{[1]} \end{bmatrix} + g, \quad \underset{q \neq q \neq q}{\text{with}} \\ \sigma(\sqrt{s_{\gamma p}}) = f_{i}(x_{1}, \mu_{F}) \otimes \hat{\sigma}(\eta), \quad \underset{q \neq q \neq q}{\text{with}} \\ \text{where } \eta = \frac{\hat{s} - M^{2}}{M^{2}} \text{ with } \hat{s} = (q + p_{1})^{2}, z = \frac{pP}{qP}. \quad \underset{q \neq q \neq q \neq q}{\text{with}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} + u_{0} u_{s} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_{0} u_{0} u_{0} u_{0}}} \\ \overset{q = (q + p_{1})^{2}}{\sqrt{s_{s} - u_{0} u_$$



3

Solution: Matching between High-Energy Factorisation and NLO

We resum corrections $\propto \alpha_s^n \ln^{n-1}(\hat{s}/M^2)$ to $\hat{\sigma}(\hat{s})$ using HEF and smoothly match it with NLO CF by the *Inverse-Error Weighting Method*.



3

 η_c -hadroproduction [Lansberg, Nefedov, Ozcelik 21']

Collaborative activities

- Visit of A. van Hameren (IFJ PAN, Krakov) to IJClab, Orsay on 14/08-18/08/2023, funded by the Marie Sklodowska Curie action "RadCor4HEF" (grant agreement t No. 101065263)
- Participation of A. van Hameren (IFJ PAN, Krakov), K. Kutak (IFJ PAN, Krakow), G. Beuf (NCBJ, Warsaw) and T. Altinoluk (NCBJ, Warsaw) in the "Mini-Workshop on the overlap between QCD resummations" in the Centre Paul Langevin (CNRS) in Aussois (France) on 14th 17th of January 2024. Funded by COPIN.

Transversity GPDs (chiral-odd GPDs) are completely unknown experimentally. We have proposed the exclusive photoproduction of a photon-meson pair with large invariant mass to extract them:

 $\blacktriangleright \gamma N \to \gamma M N':$

R. Boussarie, G. Duplančić, S. Nabeebaccus, K. Passek-Kumerički, B. Pire,

L. Szymanowski, S. Wallon: [1609.03830, 1809.08104, 2212.00655, 2302.12026]

Moreover, the richer kinematics of the process allows the sensitivity of GPDs wrt x to be probed (beyond moment-type dependence, e.g. in DVCS): J. Qiu, Z. Yu: [2305.15397]

Exclusive photoproduction of γ -meson pair with large invariant mass

► Consider the process $\gamma N \rightarrow \gamma M N'$, M =meson. Collinear factorisation of the amplitude at large $M^2_{\gamma M}$, t', u', and small t.



Factorised formula:

$$\mathcal{A} = \int_{-1}^{1} dx \int_{0}^{1} dz \ T(x,\xi,z) \ H(x,\xi,t) \ \Phi_{M}(z)$$

 $T(x, \xi, z)$: Coefficient function (Hard), computed in pQCD $H(x, \xi, t)$: GPD (soft) $\Phi_M(z)$: Distribution amplitude of outgoing meson

- Mesons considered in the final state: π^{\pm} , $\rho_{L,T}^{\pm,0}$.
- Leading order and leading twist.
- Cross-section (and polarisation asymmetries wrt to incoming photon) differential in (−u') and M²_{γM}, and evaluated at (−t) = (−t)_{min}, covering S_{γN} from ~ 4 GeV² to 20000 GeV².
- Very good statistics at various experiments, such as JLab, COMPASS, future EIC and LHC in ultra-peripheral collisions (UPCs).
- Small ξ studies can also be done using experimental data from collider experiments.

Extension of the calculation to NLO $_{\mbox{Quark GPD case}}$

At LO, there are 20 diagrams, but at NLO, this goes up to 422!

 \implies Necessary to automate!

Our approach:

- 1. Generate diagrams using FeynArts
- Reduce tensor loop integrals (which can go up to 6-point functions!) to a basis of *known scalar master integrals*.
 ⇒ Use ROLI (Reduction Of Loop Integrals), a private code based on Integration-By-Parts (IBP) reduction developed by Goran Duplancic, which is based on B. Nizic, G. Duplancic [hep-ph/0303184].
- Include GPD evolution and observe explicitly the cancellation of IR divergences.
- 4. Perform convolution over momentum fractions x (GPD) and z (DA).

Computation very similar to B. Nizic, G. Duplancic [hep-ph/0607069] for $\gamma\gamma\to\pi^+\pi^-$... except ...

Extension of the calculation to NLO Quark GPD case

- ▶ No *i* ϵ factors needed when calculating the convolution of coefficient function with 2 DAs in the $\gamma\gamma \rightarrow \pi^+\pi^-$ case.
- In γN → γMN, since poles of propagators are crossed during the convolution, one requires iε factors to be in place in arguments of logs and dilogs (easy), as well as in denominators (hard).
- ▶ Denominators can appear both through the IBP reduction procedure, or through the evalutation of master integrals themselves, where naive analytic continuation $(p_i^2 \rightarrow p_i^2 + i\epsilon)$ does NOT lead to the correct prescription! [B. Nizic, G. Duplancic: hep-ph/0006249]

Extension of the calculation to NLO Quark GPD case

- Finally, need to deal with numerical instabilities in convolution integral: These instabilities are present even in the $\gamma\gamma \rightarrow \pi^+\pi^-$ calculation, due to the introduction of spurious singularities that should cancel in the end...
- With *i*e in denominators, the situation becomes much more complicated.
- ► This was actually a significant bottleneck in the NLO computation of γN → γγN, performed by O. Grocholski, B. Pire, P. Sznajder, L. Szymanowski, J. Wagner [2110.00048, 2204.00396], where a finite iε was kept for the numerics.

 \implies However, calculation significantly simpler than our case, since only one convolution integration to perform, and also have up to 5-point functions to reduce.

- Unlike the channels considered before, having a π^0 meson in the final state allows the exchange of two gluons in the *t*-channel, due to its quantum number $J^{PC} = 0^{-+}$.
- Therefore, the amplitude becomes sensitive to gluon GPDs.
- Assuming collinear factorisation, we performed a calculation of the gluon GPD contribution, and found that it *diverges* at *leading order* and *leading twist*! This indicates a breakdown of collinear factorisation!
- ▶ In our very recent paper [S. Nabeebaccus, J. Schoenleber, L. Szymanowski, S. Wallon: 2311.09146], we were able to show that this is due to the existence of *Glauber gluon* $(|k_T|^2 \gg k^+k^-)$ exchanges which remain trapped, preventing a *collinear* factorisation of the process.
- However, all processes we have previously studied are safe, since two-gluon exchanges are forbidden.

Heavy-quark constraints on nuclear PDFs

- nPDFs: essential link between measurable hadronic cross-sections and calculable cross-sections induced by partons
- More challenging than nucleon PDFs (not just a sum of nucleon PDFs)
- Determined by performing global analyses of experimental data
- Gluon content of nuclei very little known compared to quark content (x < 10⁻³ region extrapolated from larger x region)
- Necessary for phenomenology of heavy-ion to determine the small-x gluon nPDFs and to reduce their uncertainties
- ▶ Recent study of the impact of heavy quark data
 ⇒ Way to constrain gluon density down to x = 7.10⁻⁶
 A. Kusina, J.P. Lansberg, I. Scheinbein, H.S. Shao, Phys.Rev.D 104 (2021) 1, 014010
 P. Duwentäster, T. Ježo, M. Klasen, K. Kovařík, A. Kusina, Phys.Rev.D 105 (2022) 11, 114043

Heavy-quark constraints on nuclear PDFs

- Focus on the spatial dependence of the gluon nPDFs
- Assumption that the spatially dependent nuclear modification can be determined by the local nucleon number density (thickness function)

$$R^{\mathcal{A}}(\mathbf{b},x,\mu_{F})-1=(R^{\mathcal{A}}(x,\mu_{F})-1)G\left(rac{T_{\mathcal{A}}(\mathbf{b})}{T_{\mathcal{A}}(0)}
ight),$$

with

$$\int T_{\mathcal{A}}(\mathbf{b}) G\left(\frac{T_{\mathcal{A}}(\mathbf{b})}{T_{\mathcal{A}}(0)}\right) d^2\mathbf{b} = \mathcal{A}.$$

We take:

$$G\left(rac{T_{\mathcal{A}}(\mathbf{b})}{T_{\mathcal{A}}(0)}
ight) \propto \left(rac{T_{\mathcal{A}}(\mathbf{b})}{T_{\mathcal{A}}(0)}
ight)^{\gamma_{\mathcal{A}}},$$

with γ_A to be determined ($\gamma_A = 1$ often used in literature)

Heavy-quark constraints on nuclear PDFs

- 2 types of data are used:
 - Centrality dependent R_{pA} for single inclusive particle production (in our case pAu $\rightarrow J/\Psi$ forward data from PHENIX (RHIC) at $\sqrt{s_{NN}} = 200 \text{ GeV}$)
 - * forward: tension between nPDFs and backwards data
 - * J/ Ψ : Comover effect may be important for excited states like $\Psi(2S)$
 - * RHIC: Large event-by-event fluctuations for LHC data.
 - Double Parton Scattering (DPS) in minimum bias pA collisions.
 - * Choice of pPb $\rightarrow D^0 D^0$ LHCb data
 - * pPb $\rightarrow J/\Psi D^0$ data also available but suffer from large SPS contamination.

Centrality dependent R_{pA}



$$P_{pAu}(b_{min} < b < b_{max}) = 1 - r + rR_{pAu}$$

$$r = \frac{\int_{b_{min}}^{b_{max}} T_A(\mathbf{b}) G\left(\frac{T_A(\mathbf{b})}{T_A(\mathbf{0})}\right) d^2\mathbf{b}}{\int_{b_{min}}^{b_{max}} T_A(\mathbf{b}) d^2\mathbf{b}}$$

- ► R_{pAu} obtained from experiment (0-100% centrality)
- Best fit gives γ_A = 1.50 ± 0.10 with χ²/ndf = 0.64.
 (24 data points vs y_{cms}(J/Ψ) and 96 vs p_T(J/Ψ))
- Still work to do (correspondence between b and the centrality classes is more complex than what is done here)

Minimum bias DPS



H.S. Shao, Phys. Rev. D 101 (2020) 5, 054036

$$\begin{array}{l} & \mathcal{R}^{DPS}_{pPb \rightarrow D_0 D_0} = \frac{\sigma^{DPS}_{pPb \rightarrow D_0 D_0}}{A \sigma^{DPS}_{pp \rightarrow D_0 D_0}} \\ = & \Sigma^2_{i,j=1} \left(\hat{\mathcal{T}}_{A,ij} + (\mathcal{A} - 1) \sigma_{eff,pp} \hat{\mathcal{T}}^{(2)}_{A,ij} \right) \times \\ & \left(\mathcal{R}^{D_0}_{pPb} \right)^{2-i} \left(\mathcal{R}^{D_0}_{pPb} \right)^{2-j} \end{array}$$

- σ_{eff,pp}: effective pp cross-section without initial parton-parton correlations, measured from DPS in pp collisions
- *R*^{D₀}_{*pPb*}: Nuclear modification of single inclusive *D*₀ production (measured)

• 2
$$\chi^2$$
 minima: $\gamma_A = 1.68 \pm 0.23$ and $\gamma_A = 2.47 \pm 0.16$

Combined results



- 2 very different observables but compatible results
- ► Atomic numbers of lead and gold are close ⇒ combination of fits
- ► Global fit gives $\gamma_A = 1.56 \pm 0.14$ (less than 10% relative uncertainty) ⇒ Highly disfavors $\gamma_A = 0$ (and also $\gamma_A = 1$ to a lesser extent)

2024 activities

- Main activities in 2024:
 - Determination of gluon nuclear PDFs (A. Kusina, J.P. Lansberg, I. Schienbein, H.S. Shao, S. Delorme): paper to be submitted
 - QCD corrections to exclusive processes to study GPDs at NLO accuracy (C. Flett, S.Nabeebaccus, J. Wagner, L. Szymanowski, S. Wallon): 1 paper on 2-to-3 processes and one on quarkonium production
 - Ph. D in direct co-supervision between France and Poland: A. Safronov, L. Manna, A. Colpani Serri
 - 4 other Ph. D from IJCLAB involved in the scientific exchanges: J. Bor, Y. Yedelkina, K. Lynch, A.C. John Rubesh Rajan

2024 activities

- The MSCF EU project of M. Nefedov on High Energy/kT factorisation for quarkonium production (in collaboration with A. van Hameren)
- ► The EU Virtual Access NLOAccess between IJCLab, LPTHE, UCLouvain and WUT, with the inclusion of TMD factorisation by C. Flett and J. Bor (PI: J.P. Lansberg)
- Review on the quarkonium physics at the EIC (editors: J.P. Lansberg, M.Nefedov, D. Kikola)